

Förra veckan

- Frekvenskurvor
- Laplacetransform - dubbelsidig och enkelsidig
- Inverkan av initialtillstånd
- Historik: Den återkopplade förstärkaren

Dagens föreläsning

- Den återkopplade förstärkaren - Nyquist
- Argumentvariationsprincipen
- Nyquistkriteriet
- Exempel
- Bodes relationer

Harry Nyquist 1889-1976

From farm life in Nilsby Värmland to Bell Labs

Dreaming to be a teacher

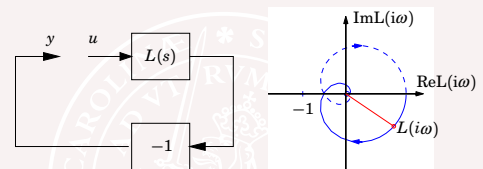
- Emigrated 1907
- High school teacher 1912
- MS EE U North Dakota 1914
- PhD Physics Yale 1917
- Bell Labs 1917

Key contributions

- Johnson-Nyquist noise
- The Nyquist frequency
- Nyquist's stability theorem



Condition for Oscillations



Cut the loop. Let u be a sinusoid. If y is a sinusoid with the same amplitude and phase, then the loop can be closed and the oscillation will be maintained. The condition for this is

$$L(i\omega) = -1$$

where $L(s) = P(s)C(s)$ is the loop transfer function. The condition implies that the Nyquist curve of $L(s)$ goes through the point -1 (the critical point)!

The Motivation 1

Mr. Black proposed a negative feedback repeater and proved by tests that it possessed the advantages predicted for it. In particular, its gain was constant to a high degree, and it was linear enough ...

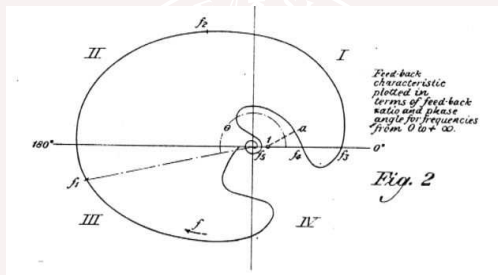
For best results, the feedback gain factor, the quantity usually known as $\mu\beta$ (the loop transfer function $L(s)$) had to be numerically much larger than unity. The possibility of stability with a feedback factor greater than unity was puzzling. Granted that the factor is negative it was not obvious how it would help. If the factor was -10 the effect of one round trip around the feedback loop is to change the original current from, say 1 to -10 . After a second trip around the loop the current becomes 100, and so forth. The totality looks much like a diverging series and it was not clear how such a succession of ever-increasing components could add to something finite and so stable as experience had shown. ...

The Motivation 2

The missing part in this argument is that the numbers that describe the successive components 1, -10 , 100, and so on, represents the steady state, whereas at any finite time many of the components have not yet reached steady state and some of them, which are destined to become very large, have barely reached perceptible magnitude. My calculations were principally concerned with replacing the infinite divergent series referred to by a series which give the actual value attained at a specific time t . The series thus obtained is convergent instead of divergent and, moreover converges to values in agreement with experimental findings.

A much simpler proof was later given based on Rouché's theorem

The Original Nyquist Curve

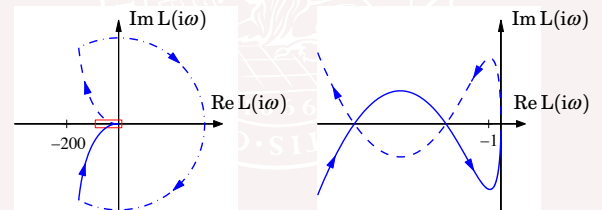


Bode moved the critical point to -1

Karl Johan Åström The Feedback Amplifier

Conditional Stability

It should perhaps be explained also how it comes to be so detailed. In the course of the calculations, the facts with which the term conditional stability have come to be associated, becomes apparent. One aspect of this is that it is possible to have a feedback loop which is stable and can be made unstable by increasing the loop gain. This seemed a very surprising result and appeared to require that all the steps be examined and set forth in full detail.



Karl Johan Åström The Feedback Amplifier

Impact of the Nyquist Theorem at ASEA

We had designed controllers by making simplified models, applying intuition and analyzing stability by solving the characteristic equation. (At that time, around 1950, solving the characteristic equation with a mechanical calculator was itself an ordeal.) If the system was unstable we were at a loss, we did not know how to modify the controller to make the system stable. **The Nyquist theorem was a revolution for us.** By drawing the Nyquist curve we got a very effective way to design the system because we know the frequency range which was critical and we got a good feel for how the controller should be modified to make the system stable. We could either add a compensator or we could use extra sensor.

Free translation from seminar by Erik Persson ABB in Lund 1970.
Why did it take 18 years?

Karl Johan Åström The Feedback Amplifier

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Argumentvariation

Låt Γ vara en enkel sluten kurva i komplexa talplanet och låt D vara det omslutna området.

Förändringen i argument för den komplexa funktionen $F(s)$ då randen till D följs moturs, kallas argumentvariationen för F längs Γ och skrivs $\Delta_{\Gamma} \arg F$:

$$\Delta_{\Gamma} \arg F := \int_{\Gamma} \left(\frac{d}{ds} \arg F(s) \right) ds$$

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Argumentvariationsprincipen

Antag att F är analytisk i en omgivning till D bortsett från ett ändligt antal poler i D . Då är

$$\frac{1}{2\pi} \Delta_{\Gamma} \arg F = N - P$$

där N är antalet nollställen och P är antalet poler för F i D .

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Bevis av argumentvariationsprincipen

Argumentet imaginärdelen av logaritmen, så

$$\begin{aligned}\Delta_{\Gamma} \arg F &= \int_{\Gamma} \left(\frac{d}{ds} \arg F(s) \right) ds \\ &= \operatorname{Im} \int_{\Gamma} \left(\frac{d}{ds} \log F(s) \right) ds = \operatorname{Im} \int_{\Gamma} \frac{F'(s)}{F(s)} ds\end{aligned}$$

F'/F är singular precis i polerna och nollställena till F .

Bevis av argumentvariationsprincipen

$$F(s) = \frac{(s - z_1) \cdots (s - z_N)}{(s - p_1) \cdots (s - p_P)} G(s)$$

where G has no poles and zeros in D . Then

$$\log F(s) = \sum_{j=1}^N \log(s - z_j) - \sum_{j=1}^P \log(s - p_j) + \log G(s)$$

Derivering och integrering ger

$$\frac{1}{2\pi} \int_{\Gamma} \frac{F'(s)}{F(s)} ds = \frac{1}{2\pi} \int_{\Gamma} \left(\sum_{j=1}^N \frac{1}{s - z_j} - \sum_{j=1}^P \frac{1}{s - p_j} + \frac{G'(s)}{G(s)} \right) ds = N - P$$

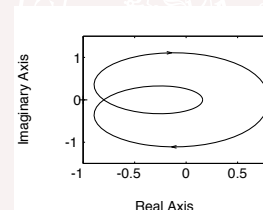
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Nyquistkriteriet

Om $L(s)$ är stabil, så är slutna systemet $[1 + L(s)]^{-1}$ stabilt om och endast om Nyquistkurvan $L(i\omega)$ inte omcirklar -1 .

Mer generellt: Skillnaden mellan antalet instabila poler i $[1 + L(s)]^{-1}$ och antalet instabila poler i $L(s)$ bestäms av antalet gånger punkten -1 omcirklas av Nyquistkurvan medurs.



Bevis av Nyquistkriteriet

Tillämpa argumentvariationsprincipen på

$$F(s) = 1 + L(s)$$

där D är det inre av en halvcirkel med centrum i origo, tillräckligt stor för att omsluta alla poler och nollställan i högra halvplanet. Då är

$$\begin{aligned}P &= \text{antalet instabila poler till } L(s) \\ N &= \text{antalet instabila poler till } [1 + L(s)]^{-1} \\ \frac{1}{2\pi} \Delta_{\Gamma} \arg F &= \text{antalet gånger punkten } -1 \text{ omcirklas av Nyquistkurvan } L(i\omega) \text{ medurs}\end{aligned}$$

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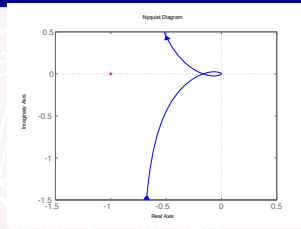
Exempel 1

$$L(s) = \frac{1}{s(s+1)(s+2)}$$

```
>> zero(1+L)
```

```
ans =
```

```
-2.3247  
-0.3376 + 0.5623i  
-0.3376 - 0.5623i
```



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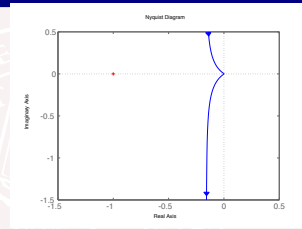
Exempel 2

$$L(s) = \frac{1}{s(s-1)(s+5)}$$

```
>> zero(1+L)
```

```
ans =
```

```
-5.0329  
0.7773  
0.2556
```



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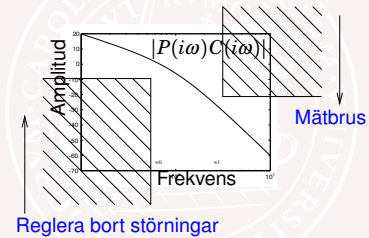
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En reglerteknisk krets ska typiskt ha hög förstärkning $|P(i\omega)C(i\omega)|$ vid låga frekvenser för att ta bort störningar och följa referenssignalen, men låg förstärkning vid höga frekvenser för att undvika stabilitetsproblem och effekterna av mätbrus.



Hur snabbt kan man gå från hög till låg förstärkning?

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Bodes relationer — Approximativ version

Om $G(s)$ är stabil och saknar nollställen i högra halvplanet så gäller

$$\arg G(i\omega_0) \approx \frac{\pi}{2} \frac{d \log |G(i\omega)|}{d \log \omega} \bigg|_{\omega=\omega_0}$$

Med nollställen i högra halvplanet blir argumentet ännu mindre.

Slutsats: Lutningen på amplitudkurvan måste vara klart mindre än 2 nära skärfrekvensen (där amplituden är ett). Annars kommer Nyquistkurvan att omcirkla -1.

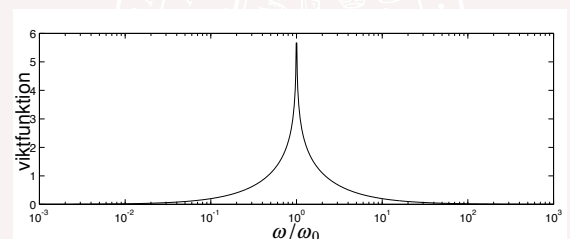
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Bode's relation — Exakt version

Om G är stabil och utan nollställen i högra halvplanet så gäller

$$\arg G(i\omega_0) = \frac{1}{\pi} \int_0^{\infty} \frac{d \log |G(i\omega)|}{d \log \omega} \underbrace{\log \left| \frac{\omega + \omega_0}{\omega - \omega_0} \right|}_{\text{viktfunktion}} d \log \omega$$



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Bodes Integralformel ("Vattensängseffekten")

For a system with loop gain $L = PC$ which has a relative degree ≥ 2 and unstable poles $p_1, \dots, p_M$, the following conservation law for the sensitivity function $S = \frac{1}{1+L}$ holds.

$$\int_0^{+\infty} \log |S(i\omega)| d\omega = \pi \sum_{i=1}^M \operatorname{Re}(p_i)$$

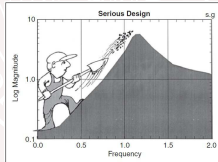


Figure 3. Sensitivity reduction at low frequency unavoidably leads to sensitivity increase at higher frequencies.

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